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X. Investigation of an Extensive Class of Partial Differential Equations of the Second Order, in which the Equation of Laplace's Functions is included.

By G. W. Hearn, Esq., of the Royal Military College, Sandhurst. Communicated by Sir John F. W. Herschel, Bart., F.R.S., &c.

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Theorem. If u be a function of x and y satisfying the equation

$$\frac{d^2u}{dxdy} + \alpha_n e^{\varphi} u = 0,$$

where

$$\frac{d^2\varphi}{dxdy} + ce^{\varphi} = 0,$$

then the solution will be

$$u = D^{-n}v_n$$

where

$$\mathbf{D} = e^{-\varphi} \frac{d}{dy},$$

where

$$v_n = \int e^{-\frac{\beta_n}{\Delta \beta_n} \phi} \chi y dy + \psi x,$$

 χy and ψx arbitrary functions of y and x,

and

$$\beta_n = \frac{\Delta \alpha_{n-1} \cdot \Delta \alpha_{n-2} \cdot \ldots}{(\Delta \alpha_{n-1} - c)(\Delta \alpha_{n-2} - c) \cdot \ldots},$$

where

$$\Delta \alpha_r = \alpha_{r+1} - \alpha_r$$

and α_n is a function of *n* vanishing for n=0 and for n=-1.

I will proceed to demonstrate this curious theorem as briefly as possible.

According to the notation, we may write the given equation

$$D\frac{du}{dx} + \alpha_n u = 0.$$

Then if

$$v_n = \mathbf{D}^n u$$

we have

$$v_{n+1}=\mathbf{D}^{n+1}u=\mathbf{D}v_n$$

$$D\frac{dv_n}{dx} + \beta_n z Dv_n = D^n \left\{ D\frac{du}{dx} + \alpha_n u \right\}. \qquad (\alpha.)$$

where z is a function of x and y to be determined, also β_n a function of u.

Writing n+1 for n in equation (α .), we ought to have

$$\mathbf{D}\frac{dv_{n+1}}{dx} + \beta_{n+1}z\mathbf{D}v_{n+1} = \mathbf{D}^{n+1}\left\{\mathbf{D}\frac{du}{dx} + \alpha_{n+1}u\right\} (\beta.)$$

This circumstance will serve to determine z and β_n as follows: we have identically

$$\mathbf{D}^{n+1} \Big\{ \mathbf{D} \frac{du}{dx} + \alpha_{n+1} u \Big\} = \mathbf{D}^{n+1} \Big\{ \mathbf{D} \frac{du}{dx} + \alpha_{n} u \Big\} + \Delta \alpha_{n} \mathbf{D}^{n+1} u$$

$$\dot{=} \mathbf{D}^{2} \frac{dv_{u}}{dx} + \beta_{n} \mathbf{D} \Big\{ z \mathbf{D} v_{n} \Big\} + \Delta \alpha_{n} \cdot \mathbf{D} v_{n} \text{ by } (\alpha)$$

$$= \mathbf{D}^{2} \frac{dv_{n}}{dx} + \beta_{n} \mathbf{D} \Big\{ z v_{n+1} \Big\} + \Delta \alpha_{n} \cdot v_{n+1}.$$
But
$$\frac{dv_{n+1}}{dx} = \frac{d}{dx} \mathbf{D} v_{n} = \frac{d}{dx} \Big(e^{-\varphi} \frac{dv_{n}}{dy} \Big)$$

$$= -\frac{d\varphi}{dx} e^{-\varphi} \frac{dv_{n}}{dy} + e^{-\varphi} \frac{d}{dy} \cdot \frac{dv_{n}}{dx}$$

$$= -\frac{d\varphi}{dx} \mathbf{D} v_{n} + \mathbf{D} \frac{dv_{n}}{dx};$$

$$\therefore \mathbf{D}^{2} \frac{dv_{n}}{dx} = \mathbf{D} \Big\{ \frac{dv_{n+1}}{dx} + \frac{d\varphi}{dx} v_{n+1} \Big\}.$$
Hence
$$\mathbf{D} \frac{dv_{n+1}}{dx} + \mathbf{D} \Big(\frac{d\varphi}{dx} v_{n+1} \Big) + \beta_{n} \mathbf{D} (z v_{n+1}) + \Delta \alpha_{n} v_{n+1}.$$

Hence

ought to be identical with $\mathbf{D} \frac{dv_{n+1}}{dx} + \beta_{n+1} z \mathbf{D} v_{n+1}$, and hence the conditions

$$\frac{d\varphi}{dx} + \beta_n z = \beta_{n+1} z$$

$$\mathbf{D} \left\{ \frac{d\varphi}{dx} + \beta_n z \right\} = -\Delta \alpha_n.$$

Eliminating z, we have

or
$$\begin{aligned} \mathrm{D}\frac{d\varphi}{dx} &= -\frac{\Delta\alpha_n\Delta\beta_n}{\beta_{n+1}} = -c, \\ \mathrm{or} & \frac{d^2\varphi}{dxdy} + ce^{\varphi} = 0, \\ \mathrm{and} & c\beta_{n+1} = \Delta\alpha_n\{\beta_{n+1} - \beta_n\}, \\ \mathrm{or} & \beta_{n+1} = \frac{\Delta\alpha_n}{\Delta\alpha_{n-c}} \cdot \beta_n; \\ & \therefore & \beta_n = \frac{\Delta\alpha_{n-1} \cdot \Delta\alpha_{n-2} \cdot \dots}{(\Delta\alpha_{n-1} - c)(\Delta\alpha_{n-2} - c) \cdot \dots} \end{aligned}$$

by these determinations we establish the formula (β) as a consequence of (α) , and therefore if the formula (α) be true for any value of n, it will be (subject to the above conditions) true for the next superior value.

Now, when n=0, $v_0=D^0u=u$, and provided α_0 and α_{-1} are each =0, α_0 and $\Delta\alpha_{-1}$ will be each 0, and α_0 and β_0 each =0, and the equation (α .) reduces to $D \frac{dv_0}{dx}$ =D $\frac{du}{dx}$, and is therefore true for u=0. Under these restrictions it will therefore be true for any positive integral value of n. Now the symbol D represents $e^{-\varphi} \frac{d}{dy}$, and therefore if U=0, $D^nU=0$, so that we have

or
$$\begin{aligned} \mathbf{D} \frac{dv_n}{dx} + \beta_n z \mathbf{D} v_n &= 0, \\ e^{-\phi} \frac{d^2 v_n}{dx dy} + \frac{\beta_n}{\Delta \beta_n} \cdot \frac{d\phi}{dx} \cdot e^{-\phi} \frac{dv_n}{dy} &= 0, \\ \\ \frac{\frac{d}{dx} \cdot \frac{dv_n}{dy}}{\frac{dv_n}{dy}} &= -\frac{\beta_n}{\Delta \beta_n} \cdot \frac{d\phi}{dx}; \end{aligned}$$

 \therefore integrating with respect to x,

$$\frac{dv_n}{dy} = e^{-\frac{\beta_n}{\Delta\beta_n}} {}^{\varphi} \chi y$$

$$v_n = \int e^{-\frac{\beta_n}{\Delta\beta_n}} {}^{\varphi} \chi y dy + \psi x,$$

$$u = D^{-n} v_n = \int e^{\varphi} \int e^{\varphi} \dots v_n dy dy \dots$$

and

the integral sign repeated n times. The theorem is therefore demonstrated.

It may be easily shown that the equation of LAPLACE's coefficients is included in the class here considered.

The equation of Laplace by a proper choice of independent variables assumes the form

$$\frac{d^2u}{dxdy} + \frac{n \cdot n + 1}{4\cos^2 \frac{y - x}{2}} \cdot u = 0.$$

Hence with reference to the preceding investigation,

Hence
$$\begin{aligned} \mathbf{D} &= \cos^2 \frac{y - x}{2} \cdot \frac{d}{dy} \text{ and } \alpha_n = \frac{n \cdot \overline{n+1}}{4}. \\ &= e^{-\varphi} = \cos^2 \frac{y - x}{2}; \\ &\therefore \qquad \frac{d\varphi}{dx} = -\tan \frac{y - x}{2} \\ &= \frac{d^2\varphi}{dx dy} + \frac{1}{2} e^{\varphi} = 0. \end{aligned}$$

Hence $c = \frac{1}{2}$ Also $\Delta \alpha_n = \frac{n+1}{2}$; $\Delta \alpha_n - c = \frac{n}{2}$;

 $\beta_n = n \text{ and } \Delta \beta_n = 1.$

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Inserting these values in the final formula, we have

$$v_n = \int \cos^{2n} \frac{y-x}{2} \chi y dy + \psi x,$$

and

$$u = \int \cos^{-2} \frac{y - x}{2} \int \cos^{-2} \frac{y - x}{2} \dots v_n dy dy \dots n \text{ times,}$$

which agrees with Mr. Hargreave's solution.